

Scattering Vector Mesons in D4/D8 model

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Outline

- 1 Motivation
 - Gauge/String duality vs. the Real QCD
 - D4/D8 pion form factors
- 2 D4/D8 Brane Model
 - Holographic QCD
- 3 Work
 - Scattering Vector Mesons
 - Wave functions
 - Quadratic radius
- 4 Summary

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From Gauge/String duality towards Real QCD

- AdS/CFT provides means to reach strong coupling gauge theory, but it is not a QCD.



To do:

Break SUSY, add quarks,
masses, \times , confinement
scale.



Holographic QCD
D4-D8 brane model

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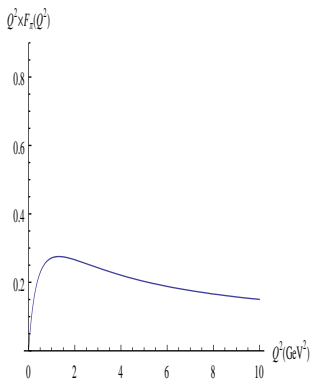
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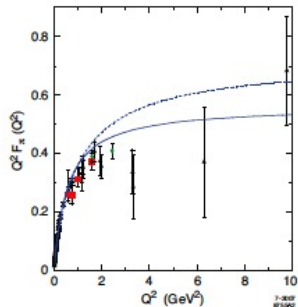
Holographic QCD
D4-D8 brane model

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(a) $Q^2 \times F_\pi(Q^2)$ plot from D4/D8 model



(b) $Q^2 \times F_\pi(Q^2)$ prediction from hardwall (dashed line), softwall (continuous line) [5] models and data

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Holographic QCD

- N_c (large) D4 branes fills $\mathcal{M}^{1,3}$, $x_4(\tau)$, with x_4 compactified in a circle.
 - 4d $SU(N_c)$ YM on the brane

- N_f probe branes D8, $\overline{D8}$ fills $\mathcal{M}^{1,3}$, $\{x_5, x_8\}$ and radial $U(\tau)$

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Holography: 4d from 5d

- quarks confined
- D8- $\overline{\text{D8}}$ gauge fields independent of $\{x_5, x_8\}$ directions.
- KK modes: vector (v_μ^n) and axial vector mesons (a_μ^n).
 - $\mathcal{A}_\mu(x^\mu, z) = \sum_{n=1}^{\infty} v_\mu^{(n)}(x^\mu) \psi_{2n-1}(z) + a_\mu^{(n)}(x^\mu) \psi_{2n}(z) + \mathcal{A}_L, \mathcal{A}_R(x)$
 ψ_n : eigenfunctions of e.o.m. (orthonormal).
 λ_n : eigenvalues $\sim m_n^2$ of (axial) vector mesons.

$$S_{D8} \sim \int d^9x e^{-\phi} \sqrt{-\det(g_{MN} + 2\pi\alpha' F_{MN})} \longrightarrow$$

$$\int d^4x \mathcal{L}_{kin+mass} - 2g_{v^n} \text{tr}(v_\mu^n \gamma^\mu) + g_{v^l v^m v^n} (\partial^\mu v^{l\nu} - \partial^\nu v^{l\mu}) [v_\mu^m, v_\nu^n]$$

$$+ g_{v^l a^m a^n} (\partial^\mu v^{l\nu} - \partial^\nu v^{l\mu}) [a_\mu^m, a_\nu^n] + g_{v^l a^m a^n} (\partial^\mu a^{n\nu} - \partial^\nu a^{n\mu}) ([v_\mu^l, a_\nu^m] - [v_\nu^l, a_\mu^m])$$

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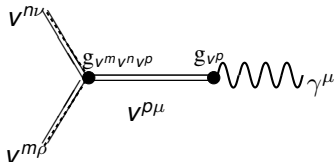


Figure: Vector meson EM scattering.

● VMD

$$\langle V^{n\nu}(k) | J^{(\mathcal{V})\mu}(0) | V^{m\rho}(k') \rangle = f_{(k,k')}^{\mu\nu\rho} F_{V^n V^m}(Q^2),$$

$$F_{V^n V^m}(Q^2) = \sum_{p=1}^{\infty} \frac{g_{V^p} g_{V^n V^m V^p}}{Q^2 + m_{V^p}^2} = \sum_{p=1}^{\infty} \frac{\frac{g_{V^p}}{m_{V^p}^2} g_{V^n V^m V^p}}{\frac{Q^2}{m_{V^p}^2} + 1} = \frac{1}{Q^2} \sum_{p=1}^{\infty} \frac{g_{V^p} g_{V^n V^m V^p}}{1 + \frac{m_{V^p}^2}{Q^2}}.$$

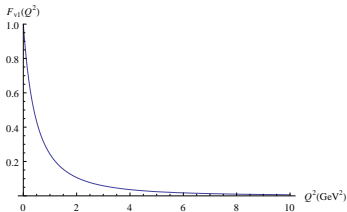
Sum Rules

Table: Masses and coupling constants

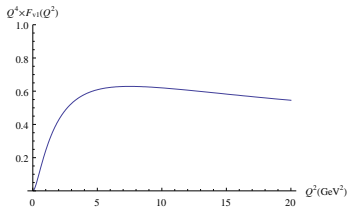
n	$\frac{m_{V^n}^2}{M_{KK}^2} = \lambda_{2n-1}$	$\kappa^{-1/2} \frac{g_{V^n}}{M_{KK}^2}$	$\kappa^{1/2} g_{V^n V^1 V^1}$
1	0.66931	2.10936	0.44658
2	2.87432	9.10785	-0.14654
3	6.59118	20.7957	1.8434×10^{-2}
4	11.79669	37.1502	-3.6885×10^{-4}
5	18.48972	58.1701	2.6953×10^{-4}
6	26.67017	83.834	3.0775×10^{-5}
7	36.33796	114.152	1.8572×10^{-5}
8	47.49318	148.103	6.9961×10^{-6}
9	53.6285	188.695	3.5081×10^{-6}

$$F_{V^m V^n}(Q^2 = 0) = \sum_{p=1}^9 \frac{g_{V^p} g_{V^n V^m V^p}}{m_{V^p}^2} \sim \delta^{mn}$$

v^1 Elastic form factor



(a) Elastic form factor F_{v^1} vs Q^2



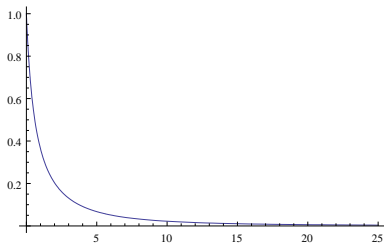
(b) $Q^4 \times F_{v^1}$ vs Q^2

Superconvergence (Grigoryan and Radyushkin):

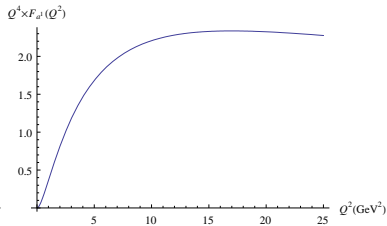
$$\sum_{n=1}^9 g_{V^n} g_{V^n v^1 v^1} \sim -0.000945 (M_{KK})^2$$

$$\sum_{n=1}^9 g_{V^n} g_{V^n a^1 a^1} \sim 0.000903 (M_{KK})^2$$

a^1 Elastic form factor



(c) Elastic form factor F_{a^1} versus Q^2



(d) $Q^4 \times F_{a^1}$ versus Q^2

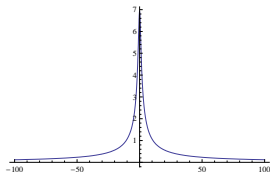
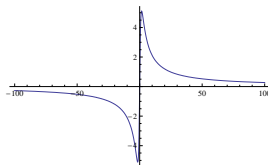
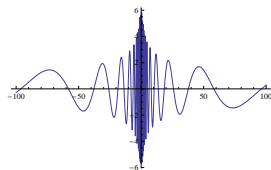
Figure: a^1 elastic form factor

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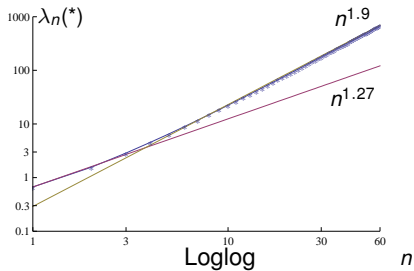
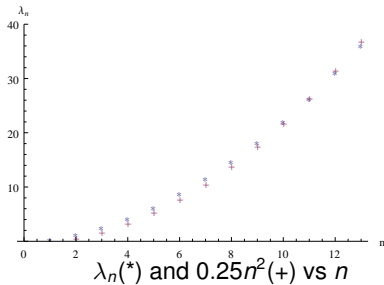
Vanishing of higher 3-vertex coupling

$$g_{V^n} = m_{V^n}^2 \kappa \int dz (1+z^2)^{-1/3} \psi_{2n-1}, \quad g_{V^1 V^1 V^p} = \kappa \int dz (1+z^2)^{-1/3} \psi_1^2 \psi_{2p-1}$$

(a) ψ_1 (b) ψ_2 (c) ψ_{60}

- ψ_1^2 selects small z region. In this region $\psi_{2p-1} \sim \cos(z\lambda_{2p-1}^{1/2})$

Regge Trajectory



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Using the ρ (vector) meson form factor expression in the low Q^2 limit, in terms of the electric form factor:

$$G_c^\rho(Q^2) = \left(1 - \frac{Q^2}{6m_{\nu_1}^2} F_{\nu_1}(Q^2) \right)$$

$$\langle r_\rho^2 \rangle = -6 \frac{d}{dQ^2} G_c^\rho(Q^2) \Big|_{Q^2=0}$$

We found






$$\langle r_\rho^2 \rangle = 0.57 \text{fm}^2$$

While results from Dyson-Schwinger equations : $\langle r_\rho^2 \rangle = 0.54 \text{fm}^2$

Summary

- Besides some weaknesses (mass spectrum, UV behavior), D4-D8 model shows adequate results, concerning vector meson EM scattering.
- Outlook
 - Would our results improve if quarks were massive?
 - What are the general features that makes D4-D8 model interesting?
 - D4-D8 is not really a Top-down model, since M_{KK} and κ are fixed by experimental values. What can be done?

For Further Reading I

-  T. Sakai and S. Sugimoto, Prog. Theor. Phys. **113** (2005) 843 [arXiv:hep-th/0412141].
-  T. Sakai and S. Sugimoto, Prog. Theor. Phys. **114** (2005) 1083 [arXiv:hep-th/0507073].
-  H. R. Grigoryan and A. V. Radyushkin, Phys. Rev. D **76**, 095007 (2007) [arXiv:0706.1543 [hep-ph]].
-  H. R. Grigoryan and A. V. Radyushkin, Phys. Lett. B **650**, 421 (2007) [arXiv:hep-ph/0703069].
-  S. J. Brodsky and G. F. de Teramond, Phys. Rev. D **77**, 056007 (2008) [arXiv:0707.3859 [hep-ph]].