

## Light-Cone 2009: Relativistic Hadronic and Particle Physics

Instituto de Física  
Universidade Federal do Rio de Janeiro

## Glueballs at finite temperature from AdS/QCD

Alex S. Miranda

Work done in collaboration with C. A. Ballon Bayona,  
Henrique Boschi-Filho, and Nelson R. F. Braga.

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# Outline

- 1 An overview on the AdS/CFT correspondence
- 2 Top-down and bottom-up holographic models
- 3 The soft-wall model at finite temperature
- 4 Thermal two-point correlation function
- 5 The spectrum of quasinormal frequencies
- 6 Concluding remarks

# 1. AN OVERVIEW ON THE ADS/CFT CORRESPONDENCE

## Original formulations [J. Maldacena, *Adv. Theor. Math. Phys.* **2**, 231 (1998)].

String theory (or M theory)  
on AdS background

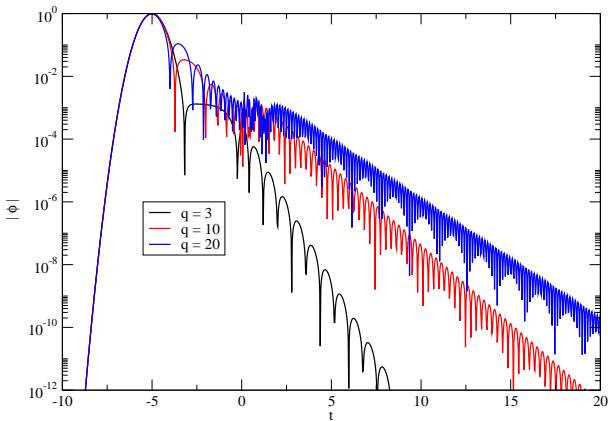
dual

Conformally invariant field  
theories (CFTs) on the boundary

- Basis for traditional applications to QCD and QGP:
  - Type IIB string theory in  $\text{AdS}_5 \times S^5$ , which is the near-horizon geometry of a set of  $N_c$  D3-branes;
  - $\mathcal{N} = 4$   $SU(N_c)$  Super-Yang-Mills theory in 3+1 dimensions.
- Basis for recent applications to condensed matter physics:
  - M-theory in the presence of  $N_c$  M2-branes, which curves the background to  $\text{AdS}_4 \times S^7$ ;
  - $\mathcal{N} = 8$   $SU(N_c)$  Super-Yang-Mills theory in 2+1 dimensions.

gravity	gauge theory
string coupling constant $4\pi g_s$	gauge coupling constant $g_{YM}^2$
dimensionless parameter $R^4/l_s^4$	't Hooft coupling constant $g_{YM}^2 N_c$
classical (super)gravity	strongly interacting theory
scalar field $\varphi$	operator $\mathcal{O} = \text{Tr } F_{\mu\nu} F^{\mu\nu}$
mass of the field	scaling dimension of the operator
isometries of the AdS space	conformal symmetries
black hole (brane)	thermal equilibrium system
parameters $(M, J, Q)$	energy, momentum, charge
$T_{Hawking}$ , $S_{Bekenstein-Hawking}$	temperature, entropy
partition function	generator of Green's functions
gravitational perturbations	fluctuations of the tensor $T_{\mu\nu}$
electromagnetic perturbations	fluctuations of the current $J_\mu$
quasinormal modes (QNMs)	poles of the correlation functions
real part of the QN frequencies	energy of the collective excitations
imaginary part of the QN freq.	thermalization time scale

# What are the black hole (brane) quasinormal modes?



## 2. TOP-DOWN AND BOTTOM-UP HOLOGRAPHIC MODELS

## Some examples of holographic models

### ■ Top-down models:

- D3/D7 branes:  $N_c$  colors and, usually,  $N_f \ll N_c$  flavors;  
A. Karch and E. Katz, *JHEP* **06**, 043 (2002).
- D4/D8 branes: flavor degrees of freedom from D8-D $\bar{8}$  pairs.  
T. Sakai and S. Sugimoto, *Prog. Theor. Phys.* **113**, 843 (2005).

### ■ Bottom-up models:

- Hard wall (HW): introduction of an IR cutoff at AdS space.  
J. J. Polchinski and M. J. Strassler, *PRL* **88**, 031601 (2002).  
H. Boschi-Filho and N.R.S. Braga, *JHEP* **05**, 009 (2003).  
G.F. de Téramond and S.J. Brodsky, *PRL* **94**, 201601 (2005).
- Soft wall (SW): presence of a dilatonlike field  $\Phi(z) = cz^2$ .  
J. A. Karch et. al., *PRD* **74**, 015005 (2006).
  - Meson spectrum:  $m_{n,S}^2 = 4c(n+S)$ ;
  - Glueball spectrum:  $m_n^2 = 4c(n+2)$ .  
J. P. Colangelo et. al., *Phys. Lett. B* **652**, 73 (2007).



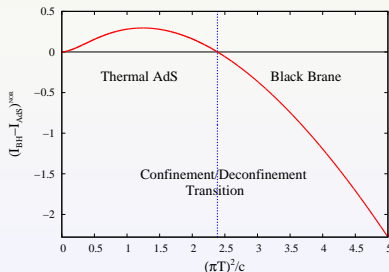
## 2. THE SOFT-WALL MODEL AT FINITE TEMPERATURE

## Action and phase transition [speed of light = $\hbar = k_B = 1$ ]

- Action for  $\varphi$  in the soft-wall model:

$$I = -\frac{\pi^3 R^5}{4\kappa_{10}^2} \int d^5x \sqrt{-g} e^{-\Phi} g^{MN} \partial_M \varphi \partial_N \varphi.$$

- Free energy in the soft-wall model: C. Herzog, *PRL* **98**, 091601 (2007); C. A. Ballon Bayona et. al., *PRD* **77**, 046002 (2008).



- Transition temperature:  $(\pi T)^2/c = 2.38643\dots$

## The black brane solution

- Spacetime metric:

$$ds^2 = e^{2A(z)} [-f(z)dt^2 + dx_1^2 + dx_2^2 + dx_3^2 + f(z)^{-1}dz^2],$$

where

$$f(z) = 1 - (z/z_0)^4$$

and

$$A(z) = -\ln(z/R).$$

- Anti-de Sitter radius:

$$R = (-6/\Lambda)^{1/2}.$$

- Hawking temperature:

$$T = 1/\pi z_0.$$

## 4. THERMAL TWO-POINT CORRELATION FUNCTION

## Massless scalar-field perturbations

- Linearized equation for  $\varphi$  on the black-brane spacetime.
- Fourier transformation of  $\varphi$  with  $k_\mu = (\omega, 0, 0, q)$ .
- In terms of  $\psi = e^{-B/2}\varphi$  ( $B = \Phi - 3A$ ), we have

$$\partial_{r_*}^2 \psi + \omega^2 \psi = V \psi,$$

where  $dz/dr_* = -f$  and  $V = f(q^2 + e^{B/2}\partial_{r_*}^2 e^{-B/2})$ .

- Normalizable and non-normalizable solutions:

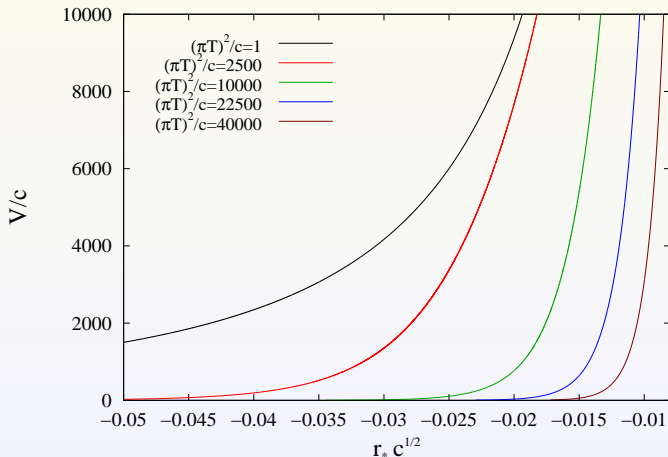
$$\psi_1(z) = z^{5/2}(1 + \dots), \quad \psi_2(z) = z^{-3/2}(1 + \dots) + a \ln(z^2) \psi_1(z).$$

- The ingoing solution at horizon  $\psi^{in}(z)$  can be written as:

$$\psi^{in}(z) = \mathcal{A}(\omega, q) \psi_2(z) + \mathcal{B}(\omega, q) \psi_1(z).$$

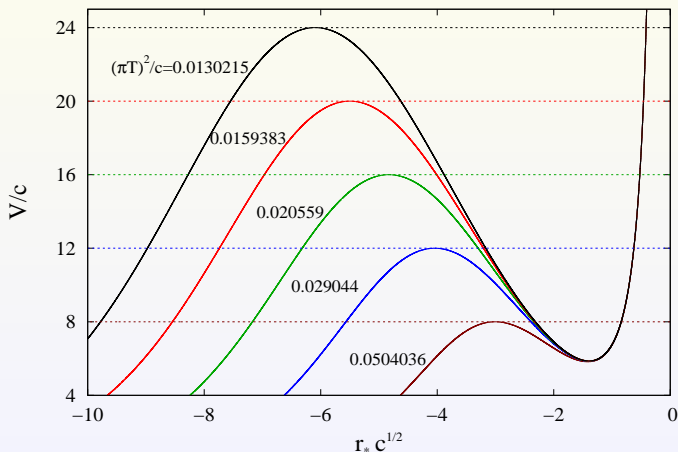
## Effective potential for high temperatures

- Potential in the high-temperature regime and  $q = 0$ .



## Effective potential for low temperatures

- Potential in the low-temperature regime and  $q = 0$ .



# The retarded Green function

- Definition:

$$G^R(\mathbf{x} - \mathbf{y}) = -i\theta(x^0 - y^0)\langle[\mathcal{O}(\mathbf{x}), \mathcal{O}(\mathbf{y})]\rangle.$$

- We study the function  $G^R$  in the momentum space.

- Real-time AdS/CFT prescription: [D. T. Son and A. O. Starinets, *JHEP* **09**, 042 (2002).]

$$G^R(k) = \frac{\pi^3 R^5}{2\kappa_{10}^2} \lim_{z \rightarrow 0} \sqrt{-g} g^{zz} e^{-\Phi} \varphi_k^*(z) \partial_z \varphi_k(z),$$

where  $\varphi(k, z) = \varphi_k(z)\varphi_0(k)$ .

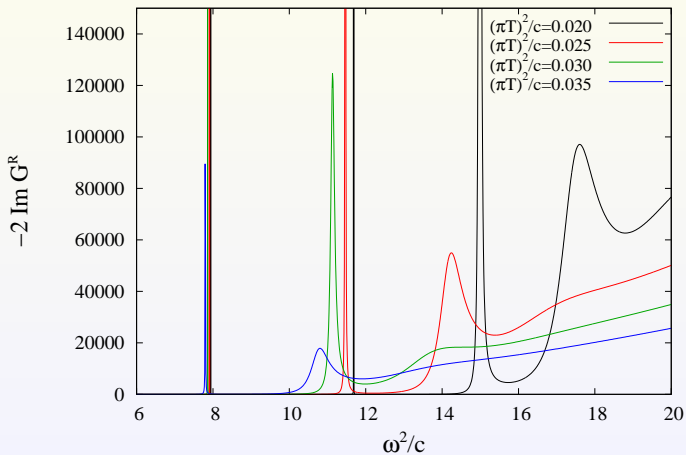
- From that follows the Green function:

$$G^R(k) = \frac{\pi^3 R^8}{2\kappa_{10}^2} \left[ 4\text{Re} \frac{\mathcal{B}}{\mathcal{A}} - i\text{Im} \frac{\mathcal{B}}{\mathcal{A}} \right] + \dots$$



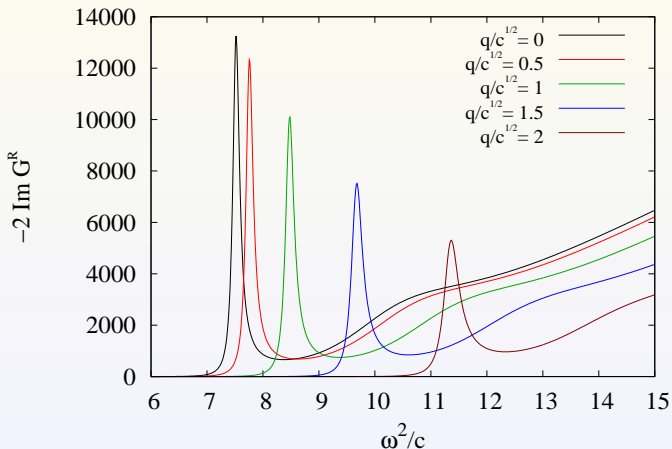
# The spectral function of scalar glueballs I

- Numerical results for  $q = 0$  and selected values of temperature



# The spectral function of scalar glueballs II

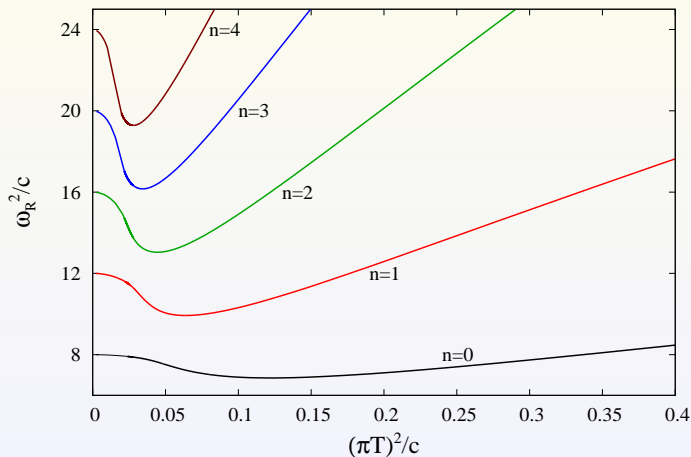
- Results for  $(\pi T)^2/c = 0.05$  and selected values of  $q/\sqrt{c}$ .



## 6. THE SPECTRUM OF QUASINORMAL FREQUENCIES

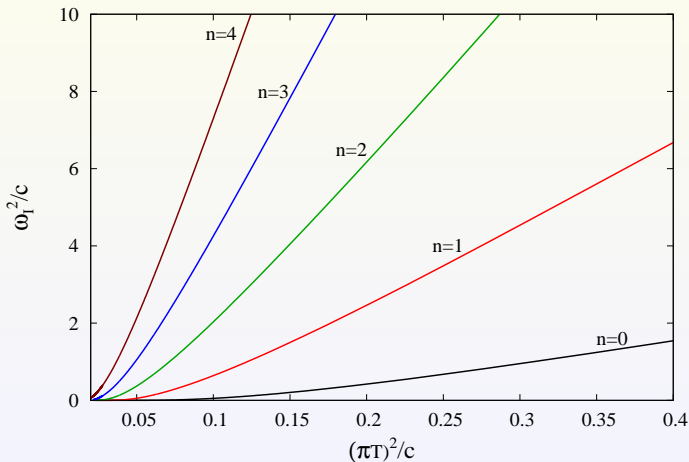
## Real part of the frequencies

- Results for the first five quasinormal frequencies and  $q = 0$



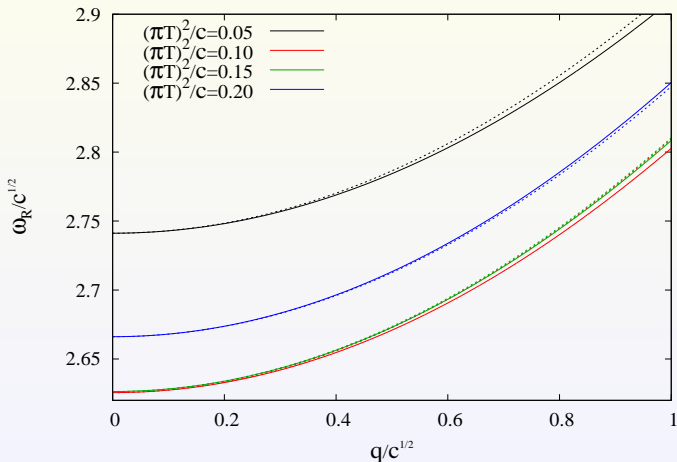
## Imaginary part of the frequencies

- Results for the first five quasinormal frequencies and  $q = 0$



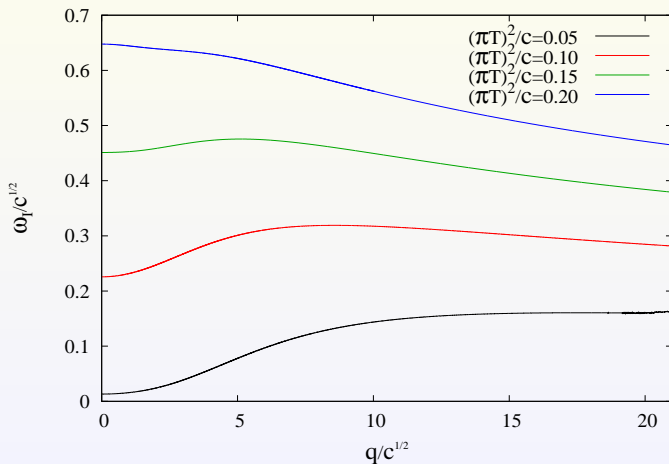
# The QNM dispersion relations

- Results for the  $n = 0$  mode and selected values of temperature



## Momentum dependence of the width

- Results for the  $n = 0$  mode and selected values of temperature



## 6. CONCLUDING REMARKS



## Overlook and perspectives

- Spectral function is featureless in the plasma phase;
- Presence of well-defined peaks in the metastable phase;
- Strong indications of glueball formation at low temperatures;
- Similar results for  $J/\psi$  mesons in the SW model;  
[M. Fujita et. al., arXiv:0903.2316 \[hep-ph\]](#)
- Similar metastable-phase behavior appeared in the D3/D7-brane model; [R. C. Myers et. al., \*JHEP\* \*\*11\*\*, 091 \(2007\)](#).
- Indication of a qualitative universal behavior!?
- More studies are necessary, including modifications of the basic SW model;  
[W. de Paula et. al., \*PRD\* \*\*79\*\*, 075019 \(2009\)](#).  
[U. Gursoy and E. Kiritsis, \*JHEP\* \*\*02\*\*, 032 \(2008\)](#).
- A derivation of the soft-wall model from string/M theory should be welcome!